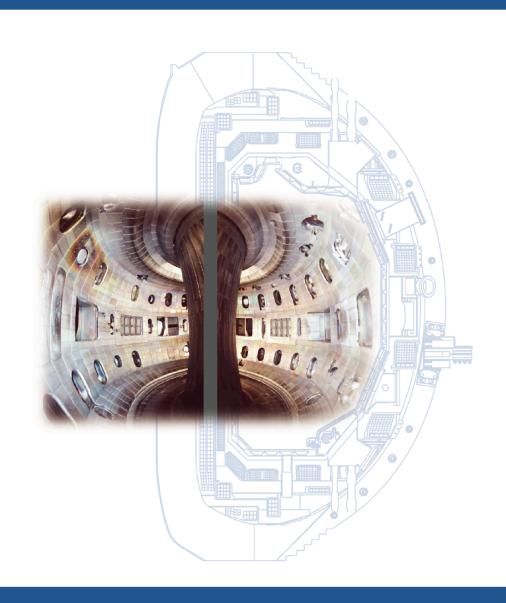
# ION ITB CHARACTERIZATION USING Te PERTURBATIONS

by

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### **MOTIVATION** and APPROACH

- Motivation: Study the characteristics of an ion Internal Transport Barrier (ITB)
  - Is the barrier of finite width or does it extend in to the plasma center?
  - How wide is the barrier if of finite width?
- Approach: Use modulated ECH to create localized heat pulses
  - have observed in the past that electron heat pulses produce a response in the ion channel
  - apply modulated ECH resonant outside the ITB
  - study ion heat pulses as they propagate through the ITB



# TARGET DISCHARGE: QH-MODE WITH A STEADY ION ITB (QDB)

 QDB has an ion ITB with characteristics most desirable for employing heat pulse propagation techniques

#### QDB characteristics

- steady state ITB allows FFT analysis of heat pulses over many modulation periods
- no ELMs or other transient events permits small perturbations to be employed ( $\delta T_e/T_e \sim 3$  5 % produces  $\delta T_i/T_i \sim 2$  3 %)

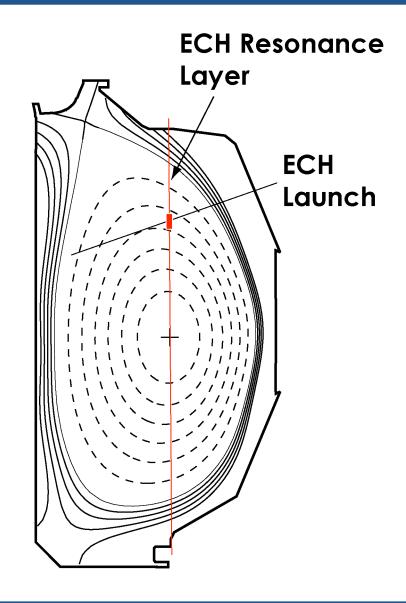
#### Focus on two cases

- Standard QDB with 4 sources of NBI,  $P_{NBI} \sim 9 MW$
- Reduce  $P_{NBI} \sim 5$  MW to try and move the location of the ITB



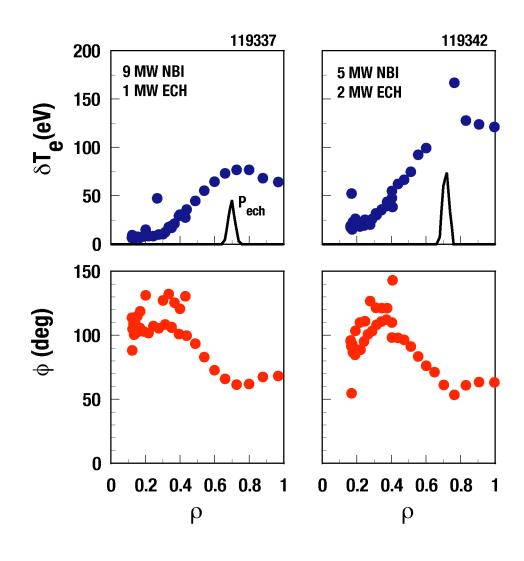
## QDB Discharge Shape Studied Is An Upper Single Null Divertor

- •1.3 MA, 2.0 T, 2 ×10<sup>19</sup> m<sup>-3</sup>
- P<sub>NBI</sub>= 5.0 and 9.2 MW
   counter injection
- Sawtooth and ELM free
  - $-q_0 = 1.1-1.2$  with monotonic profile
- ECH Resonance at  $\rho_{\text{FCH}}$ = 0.7
  - $-f_{\text{mod}} = 25 \text{ Hz}$
  - $-P_{ECH}=1-2MW$





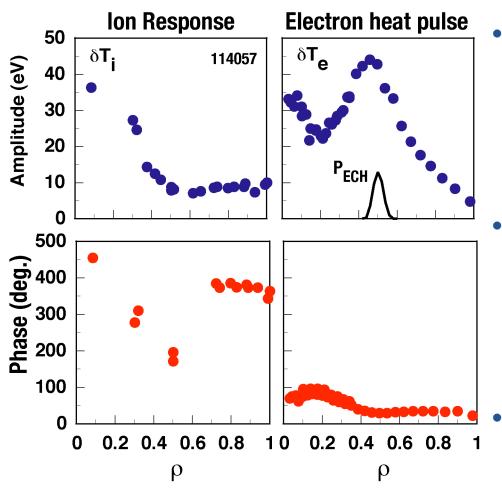
#### **ELECTRON TEMPERATURE HEAT PULSES PRODUCED BY ECH**



- T<sub>e</sub> heat pulses are characterized by their amplitude and phase
- Produced with ECH power modulated at 25 Hz and deposited at  $\rho$  = 0.7
- Small perturbations,
   δT<sub>e</sub>/T<sub>e</sub> ~ 2 3 %
- Amplitude peak and phase minimum occur at  $\rho_{\text{ECH}}$
- T<sub>e</sub> heat pulses produced a response in T<sub>i</sub>



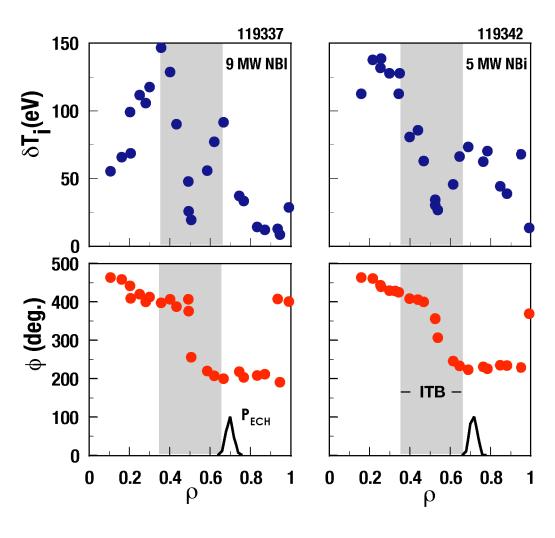
# TYPICAL T<sub>i</sub> RESPONSE TO δT<sub>e</sub> IN AN L-MODE DISCHARGE



- T<sub>i</sub> amplitude is minimum at ρ<sub>ECH</sub> and increases toward the plasma core even though δT<sub>e</sub> decreases as the pulse propagates to the plasma core ⇒ not just e-i collisional coupling
- At  $\rho_{ECH}$  the ion response is typically about 180° out of phase with the electron response ( $\phi_i \sim \phi_e$  +180), when T<sub>e</sub> increases locally, T<sub>i</sub> decreases
- This behavior is consistent with the impact of changes in  $T_e/T_i$  on ITG mode activity



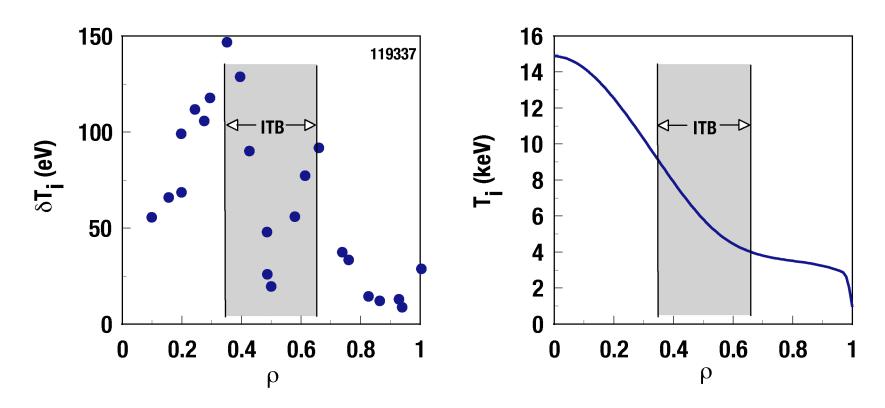
# ION RESPONSE IN QDB



- Local minimum in  $\delta T_i$  may indicate localized ion transport barrier (shaded region)
- Location of ITB is not too sensitive to ∇T<sub>i</sub> or V<sub>φ</sub> since location does not vary for 9 MW NBI compared to 5 MW NBI case



# ITB CENTERED AT $\rho$ = 0.5 AND IS ~ 15 cm WIDE



• Decrease in  $\delta T_i$  is expected in region of decreased transport. As pulse enters ITB region, propagating from  $\rho_{ECH}$ =0.7, amplitude of pulse  $\delta T_i$  decreases and trend reverses as leave ITB region



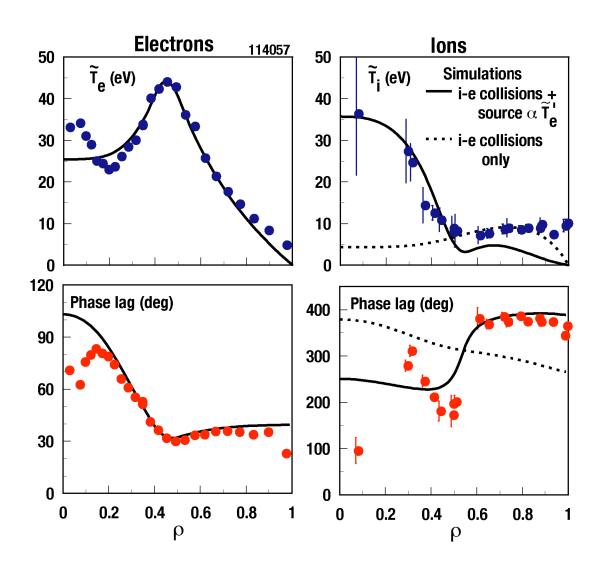
# WORK HAS BEGUN ON COUPLING ION TRANSPORT TO ELECTRONS THROUGH A DEPENDENCE OF ION TRANSPORT COEFFICIENTS ON $T_e$ AND $\nabla T_e$

- Equation for ion energy conservation
  - $3/2 \partial n_i T_i / \partial t + \nabla q_i = S_i$ where heat flux  $q_i = -n_i \chi_i \nabla T_i + n_i T_i U_i$
  - $\chi_i$  diffusion,  $U_i$  convection coefficients
- Two sources of coupling to electrons
  - source term,  $S_i = Q_{NBI} Q_{ie}$  and collisional coupling  $Q_{ie} = n^2(T_i T_e)/T_e^{3/2}$
  - ion transport coefficients may depend on  $T_e$  and/or  $\nabla T_e$  ,  $\chi_i(\rho,T_i,\nabla T_i,T_e,\nabla T_e$  ) and  $U_i(\rho,T_i,T_e)$
- Fourier transformed, linearized equation for  $\widetilde{T}_i$  with effective diffusion, convection and damping terms, D, V,  $1/\tau$ . Note there is a spatially distributed source.

$$-\mathbf{D}\widetilde{\mathsf{T}_{\mathsf{i}}}^{\prime\prime} + \mathbf{V}\widetilde{\mathsf{T}_{\mathsf{i}}}^{\prime} + (\frac{1}{\mathbf{\tau}} + i\frac{3}{2}\omega)\widetilde{\mathsf{T}_{\mathsf{i}}} = \mathsf{C}_{\mathsf{o}}\widetilde{\mathsf{T}}_{\mathsf{e}} + \mathsf{C}_{\mathsf{1}}\widetilde{\mathsf{T}}_{\mathsf{e}}^{\prime} + \mathsf{C}_{\mathsf{2}}\widetilde{\mathsf{T}}_{\mathsf{e}}^{\prime\prime}$$



## COLLISIONAL TERM ALONE CAN NOT EXPLAIN ION RESPONSE



- Collisional term alone can not account for ion perturbed response since T

  e decreases toward plasma core whereas T

  i amplitude increases.
- Not able to compute effective transport coefficients yet, but can model them with simple polynomials.
- Much better fit to ion response obtained by including a source term  $\propto \tilde{T}_e^{'}$ .



#### **SUMMARY**

- Evidence for an ion ITB has been observed in QH-mode discharges in DIII-D based on measurements of  $T_i$  perturbations driven by  $T_e$  heat pulses created with modulated ECH power.
- The ion ITB appears to be spatially localized at  $\rho$  ~ 0.5 with a width of ~ 15 cm.
- The location of the barrier and its width were insensitive to changes in  $\nabla T_i$  (factor 2.6 increase) and toroidal rotation velocity  $V_{\varphi}$  (60% increase) obtained by doubling  $P_{NBI}$ .
- Work has begun to solve the full second order partial differential equation for transport of ion perturbations with ion transport coefficients that are dependent on  $T_e$  and  $\nabla T_e$ .
  - collisional coupling alone can not account for the ion perturbations observed
  - some coupling between ion transport coefficients and electrons is required

